Robust Determination of the Higgs Couplings: Power to the Data

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January 2013

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1 / 26

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Outline





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2 The (non)-official status



Analysis Framework





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The Higgs boson ¹ has been (one of) the only missing piece of the SM for long \rightarrow **A** particle directly related to the EWSB.

The EWSB is one of the sectors with less experimental information, but a lot of interesting theoretical open questions arise from it \rightarrow motivations for NP modifying the sector

Studying the couplings of the Higgs could be one of the fastest tracks to NP

¹ or whatever...

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3 / 26

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The (non)-official status

The experimental research

After LEP and Tevatron, now it was the time for the LHC.

At the LHC the SM Higgs boson has to be searched through the following production mechanisms and decay modes:



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The experimental research

SM main discovery modes for $\simeq 125$ GeV: $pp \rightarrow \gamma\gamma$ $pp \rightarrow ZZ \rightarrow \ell\ell\ell\ell$ $pp \rightarrow WW \rightarrow \ell\nu\ell\nu$ Also may be possible in:



 $pp \rightarrow b\bar{b}$ (only VH)



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¡Eureka!



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26

CMS Status

(Last update mid-November 2012)*2

6.9 σ for 125.8 GeV (Nov 12)

 $\gamma\gamma$: "Driving" the excess, but not updated since July... (still 5.3 fb⁻¹ for 8 TeV).

Total 8 TeV run collected:

 $\mathcal{L} = 21.79 \text{ fb}^{-1}$



²In December first $Z\gamma$ results presented: with 5.0 fb⁻¹ for 7 TeV and 5.2 fb⁻¹ for 8 TeV sensitivity is about O(10) x SM

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ATLAS Status

(Last update mid-December 2012)³

 $\sim 7\sigma$ for 125 GeV (Dec 12)³

All analyses with $\sim 5 \text{ fb}^{-1}$ for 7 TeV and $\sim 13 \text{ fb}^{-1}$ for 8 TeV

Total 8 TeV run collected

 \mathcal{L} =21.7 fb⁻¹



 ^{3}ZZ and $\gamma\gamma$ mass measurement discrepance unexplained $\langle \mathcal{P} \rangle$ $\langle \mathbb{P} \rangle$ $\langle \mathbb{P} \rangle$ $\langle \mathbb{P} \rangle$

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Signal strenghts

Are the starting points of many analysis: $\mu = \frac{\sigma_{obs}}{\sigma_{SM}}$



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2012 summary / 2013 perspectives

Finally, 2012 has been the year of the well stablished discovery of the Higgs(-like) particle

We will have a more clear picture of the properties during 2013:

- $\bullet ~\sim 10~{\rm fb^{-1}}$ 8 TeV data to analyze in almost all of the channels, in some of them even more...
- Will CMS $\gamma\gamma$ new analysis dilute the excess?
- ATLAS ZZ and $\gamma\gamma$ mass discrepance still needs explanation.
- New parity, spin, $Z\gamma$ etc analyses expected.

But now let's look in detail to the current status of the study of the Higgs couplings

10 / 26

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Overview

Discovery of a ${\simeq}125$ GeV "Higgs-like" particle ${\rightarrow}{\rm EWSB}$ direct exploration:

- Spin
- Parity
- EWSB connected new states
- Couplings

Bottom-up model-independent effective Lagrangian approach:

$$\mathcal{L}_{\text{eff}} = \sum_{n} \frac{f_n}{\Lambda^2} \mathcal{O}_n$$

- Assume observed state is light electroweak doublet scalar and that $SU(3)_c \otimes SU(2)_L \otimes U(1)_Y$ is linearly realized in the effective theory.
- \mathcal{L}_{eff} : describe the low energy effects of new physics in the couplings of this observed new state in the coefficients of dimension-6 operators.
- Choice of operators and basis \rightarrow **Driven by the data**

Determine coefficients of operators using all available data: Tevatron, LHC, TGV, EWPD. イロト イクト イラト イラト ミークへの

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Effective Lagrangian

$$\mathcal{L}_{\text{eff}} = \sum_{n} \frac{f_n}{\Lambda^2} \mathcal{O}_n$$

Our assumptions are:

- The observed state belongs to a SU(2) doublet.
- The state is CP-even as in SM.
- Narrow resonance and no overlapping resonances.
- $\bullet~SU(3)_c\otimes SU(2)_L\otimes U(1)_Y$ SM local symmetry, C and P even, lepton and baryon number conservation

59 dimension-6 operators are enough...⁴

But we can use EOM to eliminate 3 of them:

$$2\mathcal{O}_{\Phi,2} - 2\mathcal{O}_{\Phi,4} = \sum_{ij} \left(y_{ij}^e \mathcal{O}_{e\Phi,ij} + y_{ij}^u \mathcal{O}_{u\Phi,ij} + y_{ij}^d (\mathcal{O}_{d\Phi,ij})^\dagger + \text{h.c.} \right) ,$$

$$2\mathcal{O}_B + \mathcal{O}_{WB} + \mathcal{O}_{BB} + {g'}^2 \left(\mathcal{O}_{\Phi,1} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{{g'}^2}{2} \sum_i (-\frac{1}{2} \mathcal{O}_{\Phi L,ii}^{(1)} + \frac{1}{6} \mathcal{O}_{\Phi Q,ii}^{(1)} - \mathcal{O}_{\Phi e,ii}^{(1)} + \frac{2}{3} \mathcal{O}_{\Phi u,ii}^{(1)} - \frac{1}{3} \mathcal{O}_{\Phi d,ii}^{(1)} \right)$$

$$-\frac{1}{3} \mathcal{O}_{\Phi d,ii}^{(1)}$$

$$2\mathcal{O}_W + \mathcal{O}_{WB} + \mathcal{O}_{WW} + g^2 \left(\mathcal{O}_{\Phi,4} - \frac{1}{2} \mathcal{O}_{\Phi,2} \right) = -\frac{g^2}{4} \sum_i \left(\mathcal{O}_{\Phi L,ii}^{(3)} + \mathcal{O}_{\Phi Q,ii}^{(3)} \right) .$$

[™]Buchmuller & Wyler; Grzadkowski et al. arXiv: 1008.4884 - ▲ ■ ▶ ▲ 중 ▶ ▲ 돌 ▶ ▲ 돌

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 12 / 26

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Higgs interactions with gauge bosons⁵:

$$\begin{split} \mathcal{O}_{GG} &= \Phi^{\dagger} \Phi \ G^{a}_{\mu\nu} G^{a\mu\nu} \ , \qquad \mathcal{O}_{WW} &= \Phi^{\dagger} \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \ , \qquad \mathcal{O}_{BB} &= \Phi^{\dagger} \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi \ , \\ \mathcal{O}_{BW} &= \Phi^{\dagger} \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \ , \qquad \mathcal{O}_{W} &= (D_{\mu} \Phi)^{\dagger} \hat{W}^{\mu\nu} (D_{\nu} \Phi) \ , \qquad \mathcal{O}_{B} &= (D_{\mu} \Phi)^{\dagger} \hat{B}^{\mu\nu} (D_{\nu} \Phi) \ , \\ \mathcal{O}_{\Phi,1} &= (D_{\mu} \Phi)^{\dagger} \Phi \ \Phi^{\dagger} (D^{\mu} \Phi) \ , \qquad \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^{\mu} \left(\Phi^{\dagger} \Phi \right) \partial_{\mu} \left(\Phi^{\dagger} \Phi \right) \ , \qquad \mathcal{O}_{\Phi,4} &= (D_{\mu} \Phi)^{\dagger} (D^{\mu} \Phi) \left(\Phi^{\dagger} \Phi \right) \ , \end{split}$$

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$$\begin{array}{ll} \mathcal{O}_{e\Phi,ij} = (\Phi^{\dagger}\Phi)(\bar{L}_{i}\Phi e_{R_{j}}) & \mathcal{O}_{\Phi L,ij}^{(1)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{L}_{i}\gamma^{\mu}L_{j}) & \mathcal{O}_{\Phi L,ij}^{(3)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{L}_{i}\gamma^{\mu}\sigma_{a}L_{j}) \\ & \bigoplus \\ \mathcal{O}_{u\Phi,ij} = (\Phi^{\dagger}\Phi)(\bar{Q}_{i}\tilde{\Phi}u_{R_{j}}) & \mathcal{O}_{\Phi Q,ij}^{(1)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{Q}_{i}\gamma^{\mu}Q_{j}) & \mathcal{O}_{\Phi Q,ij}^{(3)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{Q}_{i}\gamma^{\mu}\sigma_{a}Q_{j}) \\ & \bigoplus \\ \mathcal{O}_{d\Phi,ij} = (\Phi^{\dagger}\Phi)(\bar{Q}_{i}\Phi d_{R_{j}}) & \mathcal{O}_{\Phi e,ij}^{(1)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{e}_{R_{i}}\gamma^{\mu}e_{R_{j}}) \\ & \bigoplus \\ \mathcal{O}_{\Phi u,ij}^{(1)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{u}_{R_{i}}\gamma^{\mu}d_{R_{j}}) \\ & \bigoplus \\ \mathcal{O}_{\Phi u,ij}^{(1)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{d}_{R_{i}}\gamma^{\mu}d_{R_{j}}) \\ & \bigoplus \\ \mathcal{O}_{\Phi u,ij}^{(1)} = \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{u}_{R_{i}}\gamma^{\mu}d_{R_{j}}) \end{array}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

Z properties, W decays, low energy ν scattering, atomic P, FCNC, Moller scattering P and $e^+e^- \rightarrow f\bar{f}$ at LEP2.

Higgs interactions with gauge bosons⁵:

$$\begin{split} \mathcal{O}_{GG} &= \Phi^{\dagger} \Phi \ G^{a}_{\mu\nu} G^{a\mu\nu} \ , \qquad \mathcal{O}_{WW} &= \Phi^{\dagger} \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \ , \qquad \mathcal{O}_{BB} &= \Phi^{\dagger} \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi \ , \\ \mathcal{O}_{BW} &= \Phi^{\dagger} \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \ , \qquad \mathcal{O}_{W} &= (D_{\mu} \Phi)^{\dagger} \hat{W}^{\mu\nu} (D_{\nu} \Phi) \ , \qquad \mathcal{O}_{B} &= (D_{\mu} \Phi)^{\dagger} \hat{B}^{\mu\nu} (D_{\nu} \Phi) \ , \\ \mathcal{O}_{\Phi,1} &= (D_{\mu} \Phi)^{\dagger} \Phi \ \Phi^{\dagger} (D^{\mu} \Phi) \ , \qquad \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^{\mu} \left(\Phi^{\dagger} \Phi \right) \partial_{\mu} \left(\Phi^{\dagger} \Phi \right) \ , \qquad \mathcal{O}_{\Phi,4} &= (D_{\mu} \Phi)^{\dagger} (D^{\mu} \Phi) \left(\Phi^{\dagger} \Phi \right) \end{split}$$

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 EWPD at tree level, TGV

Higgs interactions with gauge bosons⁵:

$$\begin{split} \mathcal{O}_{GG} &= \Phi^{\dagger} \Phi \ G^{a}_{\mu\nu} G^{a\mu\nu} \ , \qquad \mathcal{O}_{WW} &= \Phi^{\dagger} \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \ , \qquad \mathcal{O}_{BB} &= \Phi^{\dagger} \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \Phi \ , \\ \mathcal{O}_{BW} &= \Phi^{\dagger} \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \Phi \ , \qquad \mathcal{O}_{W} &= (D_{\mu} \Phi)^{\dagger} \hat{W}^{\mu\nu} (D_{\nu} \Phi) \ , \qquad \mathcal{O}_{B} &= (D_{\mu} \Phi)^{\dagger} \hat{B}^{\mu\nu} (D_{\nu} \Phi) \ , \\ \mathcal{O}_{\Phi,1} &= (D_{\mu} \Phi)^{\dagger} \Phi \ \Phi^{\dagger} (D^{\mu} \Phi) \ , \qquad \mathcal{O}_{\Phi,2} &= \frac{1}{2} \partial^{\mu} \left(\Phi^{\dagger} \Phi \right) \partial_{\mu} \left(\Phi^{\dagger} \Phi \right) \ , \qquad \mathcal{O}_{\Phi,4} &= (D_{\mu} \Phi)^{\dagger} (D^{\mu} \Phi) \left(\Phi^{\dagger} \Phi \right) \end{split}$$

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$$\begin{split} \mathcal{O}_{e\Phi,33} &= (\Phi^{\dagger}\Phi)(\bar{L}_{3}\Phi e_{R_{3}}) \quad \mathcal{O}_{\Phi L,ij}^{(1)} &= \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{L}_{i}\gamma^{\mu}L_{j}) \qquad \mathcal{O}_{\Phi L,ij}^{(3)} &= \Phi^{\dagger}(\mathcal{D}_{\mu}^{a}\Phi)(\bar{L}_{i}\gamma^{\mu}\sigma_{a}L_{j}) \\ \mathcal{O}_{u\Phi,ij} &= (\Phi^{\dagger}\Phi)(\bar{Q}_{i}\tilde{\Phi}u_{R_{j}}) \qquad \mathcal{O}_{\Phi Q,ij}^{(1)} &= \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{Q}_{i}\gamma^{\mu}Q_{j}) \qquad \mathcal{O}_{\Phi Q,ij}^{(3)} &= \Phi^{\dagger}(\mathcal{D}^{a}_{\mu}\Phi)(\bar{Q}_{i}\gamma^{\mu}\sigma_{a}Q_{j}) \\ \mathcal{O}_{d\Phi,33} &= (\Phi^{\dagger}\Phi)(\bar{Q}_{3}\Phi d_{R3}) \qquad \mathcal{O}_{\Phi e,ij}^{(1)} &= \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{e}_{R_{i}}\gamma^{\mu}e_{R_{j}}) \\ &\stackrel{\leftrightarrow}{\Phi}(\bar{U}_{\mu}\phi)(\bar{u}_{R_{i}}\gamma^{\mu}u_{R_{j}}) \\ \mathcal{O}_{\Phi d,ij}^{(1)} &= \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{d}_{R_{i}}\gamma^{\mu}d_{R_{j}}) \\ \mathcal{O}_{\Phi u,ij}^{(1)} &= \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{u}_{R_{i}}\gamma^{\mu}d_{R_{j}}) \\ \mathcal{O}_{\Phi u,d,ij}^{(1)} &= \Phi^{\dagger}(\mathcal{D}_{\mu}\Phi)(\bar{u}_{R_{i}}\gamma^{\mu}d_{R_{j}}) \end{split}$$

In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

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In the absence of theoretical prejudice chose a basis where the operators are more directly related to the existing data

Z properties, W decays, low energy ν scattering, atomic P, FCNC, Moller scattering P and $e^+e^- \rightarrow f\bar{f}$ at LEP2.

$${}^{5}D_{\mu}\Phi = \left(\partial_{\mu} + i\frac{1}{2}g'B_{\mu} + ig\frac{\sigma_{a}}{2}W_{\mu}^{a}\right)\Phi, \ \hat{B}_{\mu\nu} = i\frac{g'}{2}B_{\mu\nu}, \ \hat{W}_{\mu\nu} = (i\frac{g}{2}\sigma^{a}W_{\mu\nu}^{a}) \quad \langle \Xi \rangle \quad \langle \Xi \rangle \quad \Xi \quad \Im \land \Box$$
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Effective Lagrangian for Higgs Interactions

$$\mathcal{L}_{eff} = -\frac{\alpha_s v}{8\pi} \frac{f_g}{\Lambda^2} \mathcal{O}_{GG} + \frac{f_{WW}}{\Lambda^2} \mathcal{O}_{WW} + \frac{f_W}{\Lambda^2} \mathcal{O}_W + \frac{f_B}{\Lambda^2} \mathcal{O}_B + \frac{f_{\text{bot}}}{\Lambda^2} \mathcal{O}_{d\Phi,33} + \frac{f_{\tau}}{\Lambda^2} \mathcal{O}_{e\Phi,33}$$

Unitary gauge:

$$\mathcal{L}_{\text{eff}}^{\text{HVV}} = g_{Hgg} H G^{a}_{\mu\nu} G^{a\mu\nu} + g_{H\gamma\gamma} H A_{\mu\nu} A^{\mu\nu} + g^{(1)}_{HZ\gamma} A_{\mu\nu} Z^{\mu} \partial^{\nu} H + g^{(2)}_{HZ\gamma} H A_{\mu\nu} Z^{\mu\nu}$$

+ $g^{(1)}_{HZZ} Z_{\mu\nu} Z^{\mu} \partial^{\nu} H + g^{(2)}_{HZZ} H Z_{\mu\nu} Z^{\mu\nu} + g^{(1)}_{HWW} \left(W^{+}_{\mu\nu} W^{-\mu} \partial^{\nu} H + \text{h.c.} \right)$
+ $g^{(2)}_{HWW} H W^{+}_{\mu\nu} W^{-\mu\nu}$

$$\mathcal{L}_{\rm eff}^{Hff} = g_{Hij}^f \bar{f}'_L f'_R H + \text{h.c.}$$

$$\begin{split} g_{Hgg} &= -\frac{\alpha_s}{8\pi} \frac{f_g v}{\Lambda^2} \qquad , g_{H\gamma\gamma} = -\left(\frac{g^2 v s^2}{2\Lambda^2}\right) \frac{f_{WW}}{2} \ , \\ g_{HZ\gamma}^{(1)} &= \left(\frac{g^2 v}{2\Lambda^2}\right) \frac{s(f_W - f_B)}{2c} \qquad , g_{HZ\gamma}^{(2)} = -\left(\frac{g^2 v}{2\Lambda^2}\right) scf_{WW} \ , \\ g_{HZZ}^{(1)} &= \left(\frac{g^2 v}{2\Lambda^2}\right) \frac{c^2 f_W + s^2 f_B}{2c^2} \qquad , g_{HZZ}^{(2)} = -\left(\frac{g^2 v}{2\Lambda^2}\right) \frac{c^2 f_{WW}}{2} \ , \\ g_{HWW}^{(1)} &= \left(\frac{g^2 v}{2\Lambda^2}\right) \frac{f_W}{2} \qquad , g_{HWW}^{(2)} = -\left(\frac{g^2 v}{2\Lambda^2}\right) \frac{c^2 f_{WW}}{2} \ , \\ g_{HWW}^{(1)} &= \left(\frac{g^2 v}{2\Lambda^2}\right) \frac{f_W}{2} \qquad , g_{HWW}^{(2)} = -\left(\frac{g^2 v}{2\Lambda^2}\right) f_{WW} \\ g_{Hij}^f &= -\frac{m_i^f}{v} \delta_{ij} + \frac{v^2}{\sqrt{2\Lambda^2}} f'_{f\Phi,ij} \\ &= \sqrt{2} \mathcal{R} \end{split}$$

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Effective Lagrangian for Higgs Interactions

$$\mathcal{L}_{eff} = -\frac{\alpha_{s}v}{8\pi} \frac{f_g}{\Lambda^2} \mathcal{O}_{GG} + \frac{f_{WW}}{\Lambda^2} \mathcal{O}_{WW} + \frac{f_W}{\Lambda^2} \mathcal{O}_W + \frac{f_B}{\Lambda^2} \mathcal{O}_B + \frac{f_{\text{bot}}}{\Lambda^2} \mathcal{O}_{d\Phi,33} + \frac{f_{\tau}}{\Lambda^2} \mathcal{O}_{e\Phi,33}$$

Unitary gauge:

$$\begin{aligned} \mathcal{L}_{\text{eff}}^{\text{HVV}} &= g_{Hgg} \ HG_{\mu\nu}^{a}G^{a\mu\nu} + g_{H\gamma\gamma} \ HA_{\mu\nu}A^{\mu\nu} + g_{HZ\gamma}^{(1)} \ A_{\mu\nu}Z^{\mu}\partial^{\nu}H + g_{HZ\gamma}^{(2)} \ HA_{\mu\nu}Z^{\mu\nu} \\ &+ g_{HZZ}^{(1)} \ Z_{\mu\nu}Z^{\mu}\partial^{\nu}H + g_{HZZ}^{(2)} \ HZ_{\mu\nu}Z^{\mu\nu} + g_{HWW}^{(1)} \left(W_{\mu\nu}^{+}W^{-\mu}\partial^{\nu}H + \text{h.c.}\right) \\ &+ g_{HWW}^{(2)} \ HW_{\mu\nu}^{+}W^{-\mu\nu} \end{aligned}$$

$$\mathcal{L}_{\rm eff}^{Hff} \quad = \quad g_{Hij}^f \bar{f}'_L f'_R H + {\rm h.c.}$$

$$\begin{split} g_{Hgg} &= -\frac{\alpha_s}{8\pi} \frac{f_g v}{\Lambda^2} \qquad , g_{H\gamma\gamma} = -\left(\frac{g^2 v s^2}{2\Lambda^2}\right) \frac{f_W W}{2} \ , \\ g_{HZ\gamma}^{(1)} &= \left(\frac{g^2 v}{2\Lambda^2}\right) \frac{s(f_W - f_B)}{2c} \qquad , g_{HZ\gamma}^{(2)} = -\left(\frac{g^2 v}{2\Lambda^2}\right) scf_{WW} \ , \\ g_{HZZ}^{(1)} &= \left(\frac{g^2 v}{2\Lambda^2}\right) \frac{c^2 f_W + s^2 f_B}{2c^2} \qquad , g_{HZZ}^{(2)} = -\left(\frac{g^2 v}{2\Lambda^2}\right) \frac{c^2 f_{WW}}{2} \ , \\ g_{HWW}^{(1)} &= \left(\frac{g^2 v}{2\Lambda^2}\right) \frac{f_W}{2} \qquad , g_{HWW}^{(2)} = -\left(\frac{g^2 v}{2\Lambda^2}\right) f_{WW} \end{split}$$

$$g^f_{H\,ij} = -\frac{m^j_i}{v} \delta_{ij} + \frac{v^2}{\sqrt{2}\Lambda^2} f'_{f\Phi,ij} + \frac{1}{\sqrt{2}} \delta_{ij} +$$

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January 2013

Collider data: building the χ^2

$$\chi^2 = \min_{\xi_{pull}} \sum_j \frac{(\mu_j - \mu_j^{exp})^2}{\sigma_j^2} + \sum_{pull} \left(\frac{\xi_{pull}}{\sigma_{pull}}\right)^2$$

Where

$$\mu_F = \frac{\epsilon_{gg}^F \sigma_{gg}^{ano}(1+\xi_g) + \epsilon_{VBF}^F \sigma_{VBF}^{ano} + \epsilon_{WH}^F \sigma_{WH}^{ano} + \epsilon_{ZH}^F \sigma_{ZH}^{ano} + \epsilon_{t\bar{t}H}^F \sigma_{t\bar{t}H}^{ano}}{\epsilon_{gg}^F \sigma_{gg}^{SM} + \epsilon_{VBF}^F \sigma_{VBF}^{SM} + \epsilon_{WH}^F \sigma_{WH}^{SM} + \epsilon_{ZH}^F \sigma_{ZH}^{SM} + \epsilon_{t\bar{t}H}^F \sigma_{t\bar{t}H}^{SM}} \otimes \frac{\mathrm{Br}^{ano}[h \to F]}{\mathrm{Br}^{SM}[h \to F]}$$

For the anomalous calculations:

$$\sigma_Y^{ano} = \left. \frac{\sigma_Y^{ano}}{\sigma_Y^{SM}} \right|_{tree} \left. \left. \sigma_Y^{SM} \right|_{soa} \right.$$

and

$$\Gamma^{ano}(h \to X) = \left. \frac{\Gamma^{ano}(h \to X)}{\Gamma^{SM}(h \to X)} \right|_{tree} \left. \Gamma^{SM}(h \to X) \right|_{soa}$$

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Collider data: the data points



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TGV and EWPD

Data on triple electroweak gauge boson vertices:

$$\begin{split} \mathcal{L}_{WWV} &= -ig_{WWV} \left\{ g_1^V \left(W_{\mu\nu}^+ W^{-\,\mu} V^{\nu} - W_{\mu}^+ V_{\nu} W^{-\,\mu\nu} \right) + \kappa_V W_{\mu}^+ W_{\nu}^- V^{\mu\nu} + \frac{\lambda_V}{m_W^2} W_{\mu\nu}^+ W^{-\,\nu\rho} V_{\rho}^{\,\,\mu} \right\} \\ \text{with} & \Delta g_1^Z = g_1^Z - 1 = -\frac{g^2 v^2}{8c^2 \Lambda^2} f_W \,\,, \\ & \Delta \kappa_\gamma = \kappa_\gamma - 1 = -\frac{g^2 v^2}{8\Lambda^2} \left(f_W + f_B \right) \,\,, \\ & \Delta \kappa_Z = \kappa_Z - 1 = -\frac{g^2 v^2}{8c^2 \Lambda^2} \left(c^2 f_W - s^2 f_B \right) \,\,. \end{split}$$

LEP data:

$$g_1^Z = 0.984^{+0.049}_{-0.049}$$

$$\kappa_\gamma = 1.004^{+0.024}_{-0.025}$$

with a correlation factor $\rho = 0.11$.

Data on EWPD in terms of the S,T,U parameters:

$$\begin{split} \Delta S &= 0.00 \pm 0.10 & \Delta T = 0.02 \pm 0.11 & \Delta U = 0.03 \pm 0.09 \\ \rho &= \begin{pmatrix} 1 & 0.89 & -0.55 \\ 0.89 & 1 & -0.8 \\ -0.55 & -0.8 & 1 \end{pmatrix} \end{split}$$

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S,T,U Parameters

$$\begin{split} \alpha \Delta S &= \frac{1}{6} \frac{e^2}{16\pi^2} \bigg\{ 3(f_W + f_B) \frac{m_H^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_H^2} \bigg) + \\ &+ 2 \big[(5c^2 - 2) f_W - (5c^2 - 3) f_B \big] \frac{m_Z^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_H^2} \bigg) \\ &- \big[(22c^2 - 1) f_W - (30c^2 + 1) f_B \big] \frac{m_Z^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_Z^2} \bigg) \\ &- 24c^2 f_{WW} \frac{m_Z^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_H^2} \bigg) \bigg\} , \\ \alpha \Delta T &= \frac{3}{4c^2} \frac{e^2}{16\pi^2} \bigg\{ f_B \frac{m_H^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_H^2} \bigg) \\ &+ (c^2 f_W + f_B) \frac{m_Z^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_H^2} \bigg) \\ &+ \big[2c^2 f_W + (3c^2 - 1) f_B \big] \frac{m_Z^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_Z^2} \bigg) \bigg\} , \\ \alpha \Delta U &= -\frac{1}{3} \frac{e^2 s^2}{16\pi^2} \bigg\{ (-4f_W + 5f_B) \frac{m_Z^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_H^2} \bigg) \\ &+ (2f_W - 5f_B) \frac{m_Z^2}{\Lambda^2} \log \bigg(\frac{\Lambda^2}{m_Z^2} \bigg) \bigg\} \end{split}$$

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 $\Delta \chi^2$ vrs f_X

Columns (analysis):

1st: f_g , f_{WW} , f_W , f_B 2nd: f_g , f_{WW} , f_W , f_B , f_{bot} 3rd: f_g , f_{WW} , f_W , f_B , f_{bot} , f_{τ}

Rows (parameters):

1st: f_a 2nd: f_{WW} 3rd: f_W 4th: f_B 5th: fbot

Colours/lines (data):

Solid red: Collider Dash pink: Collider + TGV Dot blue: Collider + TGV + EWPD



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BRs and production CS



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CMS vrs ATLAS



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2d correlations



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2d correlations



Image: A matrix

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Best fit and ranges

	Fit with $f_{bot} = f_{\tau} = 0$		Fit with f_{bot} and f_{τ}	
	Best fit	90% CL allowed range	Best fit	90% CL allowed range
$f_g/\Lambda^2 (\text{TeV}^{-2})$	1.3, 21.4	$[-1.2, 3.5] \cup [19, 24]$	1.3, 21.4	$[-21, 4.8] \cup [18, 44]$
$f_{WW}/\Lambda^2 (\text{TeV}^{-2})$	-0.43	$[-0.8, -0.1] \cup [2.85, 3.55]$	-0.39	$[-0.8,0] \cup [2.85,3.65]$
$f_W/\Lambda^2 (\text{TeV}^{-2})$	1.43	[-7.0, 10]	0.42	[-7.4, 7.6]
$f_B/\Lambda^2 (\text{TeV}^{-2})$	-8.4	[-30, 13]	0.42	[-7.4, 7.6]
$f_{bot}/\Lambda^2 (\text{TeV}^{-2})$			0.00, 0.90	$[-1.2, 0.20] \cup [0.70, 2.1]$
$f_{\tau}/\Lambda^2 (\text{TeV}^{-2})$			$0.02, \ 0.32$	$[-0.07, 0.13] \cup [0.2, 0.40]$
$BR^{ano}_{\gamma\gamma}/BR^{SM}_{\gamma\gamma}$	1.75	[1.15, 2.62]	1.70	[0.20, 3.00]
$BR_{WW}^{ano}/BR_{WW}^{SM}$	0.97	[0.75, 1.14]	1.02	[0.11, 1.94]
$BR^{ano}_{ZZ}/BR^{SM}_{ZZ}$	1.13	[0.78, 1.45]	1.03	[0.11, 1.96]
$BR_{bb}^{ano}/BR_{bb}^{SM}$	1.01	[0.84, 1.06]	1.04	[0.53, 1.53]
$BR_{\tau\tau}^{ano}/BR_{\tau\tau}^{SM}$	1.01	[0.84, 1.06]	0.85	[0.05, 2.25]
$\sigma_{gg}^{ano}/\sigma_{gg}^{SM}$	0.79	[0.47, 1.23]	0.79	[0.35, 8]
$\sigma_{VBF}^{ano}/\sigma_{VBF}^{SM}$	1.02	[0.92, 1.21]	1.00	[0.91, 1.13]
$\sigma_{VH}^{ano}/\sigma_{VH}^{SM}$	0.98	[0.58, 1.40]	1.02	[0.57, 1.49]

Best fit values and 90% CL allowed ranges for the combination of all available Tevatron and LHC Higgs data as well as TGV.

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Discussion and Conclusions

• Model independent analysis where the effects of new physics in the Higgs couplings are parametrized in \mathcal{L}_{eff} . $SU(2)_L$ doublet $\rightarrow SU(2)_L \times U(1)_Y$ gauge symmetry linearly realized:

$$\mathcal{L}_{\text{eff}} = \sum_{n} \frac{f_n}{\Lambda^2} \mathcal{O}_n \;\; ,$$

Choice of basis:

Power to the data \rightarrow operators whose coefficients are more easily related to existing data

$$\mathcal{L}_{eff} = -\frac{\alpha_s v}{8\pi} \frac{f_g}{\Lambda^2} \mathcal{O}_{GG} + \frac{f_{WW}}{\Lambda^2} \mathcal{O}_{WW} + \frac{f_W}{\Lambda^2} \mathcal{O}_W + \frac{f_B}{\Lambda^2} \mathcal{O}_B + \frac{f_{\text{bot}}}{\Lambda^2} \mathcal{O}_{d\Phi,33} + \frac{f_\tau}{\Lambda^2} \mathcal{O}_{e\Phi,33} \quad .$$

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Discussion and Conclusions

• Present status of the analysis using Tevatron, LHC, TGV and EWPD data:



SM compatible with data (60%-90% CL). Preference for larger-than-SM BR to photons and a smaller-than-SM gluon fusion production CS and decay BR into $\tau's$

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Discussion and Conclusions

THANK YOU!

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