Bringing Dark Matter into Focus

Mariangela Lisanti Princeton University

with Joseph Formaggio, Samuel Lee, Annika Peter, Benjamin Safdi, Siddharth Sharma, and Joshua Spitz

Direct Detection

Dark matter scatters off of nuclei in detectors

Measure recoil energy of nuclei



Direct Detection

Several different strategies for detecting recoil energy



From: Véronique Sanglard (La Thuille 2005)

Scattering Cross Section

Example: spin-independent interaction due to Higgs exchange



$$\sigma_{\chi N \to \chi N} \simeq \frac{\lambda_{\chi}^2 \lambda_q^2}{4m_h^4} \cdot \mu_{\chi N}^2$$
$$\sim (7 \times 10^{-44} \text{ cm}^2) \cdot \lambda_{\chi}^2$$

for Xe target, 125 GeV Higgs, 100 GeV DM

Spin-Independent Limit



LUX Collaboration [1310.8214]

Spin-Independent Limit



LUX Collaboration [1310.8214]



Direct Detection Signals

Total Rate

Annual Modulation

Gravitational Focusing

Scattering Rate For typical spin-independent and -dependent interactions, the differential scattering rate is given by Lab-frame velocity distribution Local DM density $\frac{dR}{dE_R} = n_{\rm dm} \left\langle v \frac{d\sigma}{dE_R} \right\rangle \propto \rho \int_{v}^{\infty} \frac{f(\mathbf{v}, t)}{v} d^3 v$ Minimum speed to induce a recoil with energy E_{nr}

A Spectrum of Possibilities

Smooth Halo

Streams

Fully Virialized <

Not Virialized

>

Smooth Halo

15 JUNE 1986 VOLUME 33, NUMBER 12 PHYSICAL REVIEW D Detecting cold dark-matter candidates Andrzej K. Drukier Max-Planck-Institut für Physik und Astrophysik, 8046 Garching, West Germany and Department of Astronomy, Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, Massachusetts 02138 Katherine Freese and David N. Spergel Department of Astronomy, Harvard-Smithsonian Center for Astrophysics, 60 Garden Street, Cambridge, Massachusetts 02138 (Received 2 August 1985) Proposed a model for the velocity distribution of dark matter Flat rotation curves imply that density falls off as $1/r^2$ Equilibrium + $\rho \sim r^{-2}$ = Maxwell-Boltzmann Isotropy ╋

Recoil Spectrum

Average over all possible DM velocities in the galactic halo





Streams

Trailing tidal debris of infalling subhalos results in 'streams' of dark matter Streams are dynamically cold, have 1D morphology





Direct Detection Signals

Total Rate

Annual Modulation

Gravitational Focusing

Annual Modulation

Dark matter signal modulates annually due to Earth's orbit about the Sun



Drukier, Freese, and Spergel [PRD] Review: Freese, ML, Savage [1209.3339]

Annual Modulation

More high-velocity particles in the summer But, scattering cross section is enhanced for low-velocity particles

$$\frac{dR}{dE_{\rm nr}} \propto \rho \int_{v_{\rm min}}^\infty \frac{f({\bf v},t)}{v} d^3 v$$

High-energy scattering events have a maximum ~June 1 Low-energy events have a maximum ~Dec 1







Modulation Spectrum

Shape of the modulation also depends on particle properties





Modulation Spectrum

Shape of the modulation also depends on particle properties





Modulation Spectrum

Shape of the modulation also depends on particle properties





Lee, ML, and Safdi [1307.5323]

Expand differential scattering rate in terms of Fourier components

$$\frac{dR}{dE_{\rm nr}} = A_0 + \sum_{n=1}^{\infty} \left[A_n \cos n\omega (t - t_0) + B_n \sin n\omega (t - t_0) \right]$$

Higher Fourier modes are enhanced for

high v_{min} scenarios (*i.e.*, light DM) local DM substructure in the halo (*i.e.*, streams)

Lee, ML, and Safdi [1307.5323]

$$E(A_1)/E(A_n) =$$
Exposure needed to observe A₁ to 95% confidence, relative to that for A_n



Lee, ML, and Safdi [1307.5323]

$$E(A_1)/E(A_n) = {\begin{array}{c} \operatorname{Exposure} \\ \operatorname{confiden} \end{array}}$$

Exposure needed to observe A_1 to 95% confidence, relative to that for A_n



Lee, ML, and Safdi [1307.5323]

$$E(A_1)/E(A_n) = {\operatorname{Expos} \atop \operatorname{confid}}$$

Exposure needed to observe A_1 to 95% confidence, relative to that for A_n





Modulation Anomalies

DAMA, 9.3σ

CoGeNT, 2.2o

NaI(Tl) target, 14 years of data

Ge target, 3.4 years of data



Direct Detection Signals

Total Rate

Annual Modulation

Gravitational Focusing

Gravitational Focusing

Lee, ML, Peter, and Safdi [1308.1953]

Sun's potential deflects incoming, unbound dark matter particles

Focusing is strongest during the Spring



K. Griest, PRD 1988. Alenazi and Gondolo [astro-ph/0608390]

Modulation Phase

Lee, ML, Peter, and Safdi [1308.1953]

Earth's orbit causes $\sim 3\%$ modulation that is extremized \sim June 1

Focusing causes ~1.5% modulation that is peaked ~March 1



A competition between two different modulation effects









Experimental Implications

Gravitational focusing results in a dependence of phase on recoil energy bin Powerful way to distinguish signal from background

> Example: 50 GeV DM, Ge target $E_{nr} \sim 40 - 41 \text{ keV}$ $E_{nr} \sim 10 - 11 \text{ keV}$ $E_{nr} \sim 2 - 3 \text{ keV}$



Example: Ge Target

For current thresholds, phase shift particularly significant for masses greater than ~15 GeV

Current advances in low-threshold technology could make shift relevant for ~8 GeV dark matter



Scattering rate in finite energy bin:

$$\bar{R}(E_{\min}, E_{\max}) = \int_{E_{\min}}^{E_{\max}} dE_{\mathrm{nr}} \frac{dR}{dE_{\mathrm{nr}}}$$

 $ar{t}_0$ is the time of maximal $ar{R}$

DAMA Revisited

DAMA signal can correspond to ~11 or 76 GeV dark matter

Both masses are in tension with null results from other experiments



DAMA Revisited

11 GeV scenario should still peak ~June 1 in each bin

Future NaI experiments (*i.e.*, SABRE) might push the threshold lower to where a phase shift would be measurable



DAMA Revisited

76 GeV scenario is affected by gravitational focusing The phase shift can be as much as a ~month in the low-energy bins



Dark Matter Disk

May form from the merger of subhalos that are dragged into the baryonic disk and disrupted

Corotates with the Galactic disk, but with a lag speed ~50 km/s



http://www.sciencedaily.com/releases/2008/09/080915210506.htm

Measuring the Dark Disk

Lee, ML, Safdi and Sharma [in progress]

Particles in the dark disk have low velocities in the lab frame

Therefore, significantly affected by gravitational focusing



Teaser

Relic Neutrinos

The "holy grail" of neutrino physics



http://www.aspera-eu.org/images/stories/files/Roadmap.pdf

Neutrino Capture

Neutrino capture on β -decaying nuclei provides a clear path forward

$$\nu_e + {}^3\mathrm{H} \rightarrow {}^3\mathrm{He} + e^-$$

No threshold on incoming neutrino energy





PTOLEMY Prototype

Beta *et al.*, arXiv: 1307/4738. Weinberg, Phys. Rev **128**, 1457 (1962).

The Case for Neutrinos

The neutrino capture rate on a single nucleus is given by

$$\lambda_{\nu} \propto \rho \int \sigma_{\rm NCB} v_{\nu} f_{\oplus}(\mathbf{v}, t) d^3 v$$

 $\sigma_{\rm NCB} v_{\nu}$ is velocity-independent to high accuracy at low neutrino energy

$$\lambda_{\nu} \propto \rho \; (\sigma_{\rm NCB} \, v_{\nu}) \left[\int f_{\oplus}(\mathbf{v}, t) \, d^3 v \right]$$

Integrates to unity!

However, the density still modulates due to gravitational focusing!

Relic Neutrino Modulation

Safdi, ML, Spitz, and Formaggio [1404.0680]

Gravitational focusing is the *only* source of modulation for relic neutrinos

Modulation fractions ~0.1-1% depending on v mass and velocity distribution

Requires a 10 kg-sized tritium target, which is feasible with PTOLEMY

Summary

Particle and astrophysics assumptions about dark matter can enhance higher-frequency modes of modulation spectrum

Unbound dark matter particles focused by Sun's gravitational potential, affecting the modulation phase

Phase shift most relevant for low-speed particles

i.e., masses greater than ~15 GeV, or lighter mass particles at low-threshold experiments